A branching particle system as a model of FKPP fronts

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GOAL: Unravel the complex genetic patterns that surround us



MOLECULAR CLOCK ASSUMPTION: Genealogical structure \Leftrightarrow Genetic structure

Genetics in exchangeable (or neutral, unstructured) populations



Universality class of genealogies: exchangeable coalescents



Rich mathematical structure

- $\boldsymbol{\mathsf{x}}$ Characterisation
 - (Pitman, Sagitov '99)
- × Statistics (e.g. SFS)
- × Data
 - (e.g. Steinruecken et al. '13)

Genetics in exchangeable (or neutral, unstructured) populations



- Universality class of genealogies: exchangeable coalescents



Genetics in exchangeable (or neutral, unstructured) populations

equal chance of reproductive success

structured populations each type has a different ability to

reproduce: **structured coal.** ?



Universality class of genealogies: exchangeable coalescents



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Macroscopic dynamic

local regulation

 $\partial_t u = \frac{1}{2} \partial_{xx} u + \frac{1}{2} u (1-u) (1+Bu)$ cooperation

B > 0 strength of cooperation

u(t,x) = density of individuals at x at time t > 0

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Travelling wave solutions

constant profile ϕ travelling at constant speed $v \equiv v(B)$ $u(t,x) = \phi(x - vt)$



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Two types of waves Stokes '76





pulled

B

speed of the linearised equation $(u \approx 0)$ **pushed**

faster than the linearised equation

$$v(B) = \begin{cases} 1 & B \in [0,2] \\ \frac{1}{2}\sqrt{\frac{B}{2}}\left(1 + \frac{2}{B}\right) & B \ge 2 \end{cases}$$

Macroscopic dynamic

local regulation

$$\partial_t u = \frac{1}{2} \partial_{xx} u + \frac{1}{2} u (1-u) (1+Bu) + \sqrt{\frac{u(1-u)}{N}} W$$

cooperation

B > 0 strength of cooperation

u(t,x) =density of individuals at x at time t > 0

The case B = 0 (pulled waves)

physical (1) and **evolutionary** properties (2) are prescribed by the dynamics at the very tip of the front. They are '**pulled** by the individuals at the tip'.

simulations, conjectures and heuristics Brunet et al. '97, '06, '06, '07, '08

(1) Fluctuations of the invasion front

 $v - v_N \sim \frac{c}{\log(N)^2}$

Mueller, Mytnik, Quastel '10 Bérard, Gouéré '12 Maillard '16 (2) Genealogy = Bolthausen-Sznitman
proved for a linear toy model
(no interaction)
Berestycki, Berestycki, Schweinsberg, '13



Macroscopic dynamic

local regulation

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u(t,x) = density of individuals at x at time t > 0

Recent breakthrough: conjectures by Birzu, Hallatschek, Korolev '18





 $\cdots B = 0$ - - B = 2

-B = 5

 $u(t,x) = \phi(x - \mathbf{v}t)$



Dyadic branching Brownian motion with branching rate $r(x) = \frac{1}{2}[(b-1)\mathbf{1}_{x<1} + 1]$ killing at **0** "critical" drift $-\mu(b)$

A model for travelling fronts?

Extension of The genealogy of branching Brownian motion with absorption, Berestycki, Berestycki, Schweinsberg 2013 ($b = 1, \mu = 1$)



Critical drift

Spectral decomposition of the critical operator $Au = \frac{1}{2}u'' - \mu u' + r(x)u$ (λ_1, v_1) principal eigenv. **Many-to-one lemma** $p_t(x, y) \approx ce^{\lambda_1 t} e^{(\frac{\mu^2}{2} - \frac{1}{2})t} e^{\mu x} v_1(x) e^{-\mu y} v_1(y) \Rightarrow \mu = \sqrt{1 + 2\lambda_1}$



Critical drift

Spectral decomposition of the critical operator $\mathcal{A}u = \frac{1}{2}u'' - \mu u' + r(x)u$ (λ_1, v_1) principal eigenv. Many-to-one lemma $p_t(x, y) \approx ce^{\lambda_1 t} e^{(\frac{\mu^2}{2} - \frac{1}{2})t} e^{\mu x} v_1(x) e^{-\mu y} v_1(y) \Rightarrow \mu = \sqrt{1 + 2\lambda_1}$ pulled/pushed regime: $\exists b_1 > 1$ such that

• if $b \in [1, b_1]$, $\lambda_1 = 0 \Longrightarrow \mu = 1$,

• if
$$b > b_1$$
, $\lambda_1 > 0$ and $\lambda_1 \nearrow b \Rightarrow \mu > 1$, $\mu \nearrow b$
 \rightarrow in this case $v_1(x) \propto e^{-\sqrt{\mu^2 - 1x}}$, $x \ge 1$



Spine = immortal particle $\Leftrightarrow h$ -transform

 $\mathcal{G}u = \frac{1}{2}\partial_{xx}u - \frac{v_1'(x)}{v_1(x)}\partial_x u$ stationary measure $\Pi(x) = v_1(x)^2 = h(x)\tilde{h}(y)$ $\mu > 1 \text{ (pushed)} \Rightarrow \Pi(x) \propto e^{-2\sqrt{\mu^2 - 1}x}, x \ge 1$

 $\boldsymbol{\rightarrow}$ the **spine** is in the **bulk**

Spine = immortal particle $\Leftrightarrow h$ -transform

 $\mathcal{G}u = \frac{1}{2}\partial_{xx}u - \frac{v_1'(x)}{v_1(x)}\partial_x u$ stationary measure $\Pi(x) = v_1(x)^2 = h(x)\tilde{h}(y)$

Let
$$\alpha := \frac{\mu + \sqrt{\mu^2 - 1}}{\mu - \sqrt{\mu^2 - 1}}$$

pulled/semipushed/fully pushed regime: $\exists b_2 > b_1$ such that

• for $b \in [1, b_1]$, $\mu = 1 \rightarrow \alpha = 1$,

• for
$$b > b_1$$
, $\alpha \nearrow b$, $\alpha(b_1) = 1$ and $\alpha(b_2) = 2$.

 $\mu > 1 \text{ (pushed)} \Rightarrow \Pi(x) \propto e^{-2\sqrt{\mu^2 - 1}x}, x \ge 1$

 \rightsquigarrow the **spine** is in the **bulk**



The semipushed regime $b \in (b_1, b_2)$



A branching particle system as a model of semi pushed fronts. T. 2024. Convergence of spatial branching processes to alpha-stable CSBPs: Genealogy of semi-pushed fronts. Foutel-Rodier, Schertzer and T. 2025+.

The semipushed regime $b \in (b_1, b_2)$



- The invasion is driven by **excursions** of particles
- These excursions generate large subfamilies in the associated genealogy, diversity loss
- Particles at the tip have a large reproductive value \Rightarrow large reproductive variance (in some sense)
- fast relaxation \Rightarrow collapse of structure

The fully pushed regime $b > b_2$



 Z_t = number of individuals $N \rightarrow \infty$

_SCALING LIMIT____

 $(\frac{1}{N}Z_{tN}) \Rightarrow (X_t) X$ is a Feller diffusion

GENEALOGY

The genealogy of the BBM converges to the Kingman coalescent

No excursion ! → "CLT"/mean field

Spectral analysis and k-spine decomposition of inhomogeneous branching Brownian motions. Genealogies in fully pushed fronts. Schertzer and T. 2024.

The fully pushed regime $b > b_2$



- The invasion is pushed by the **growth** in the front
- This generates only binary mergers in the associated genealogy, wide genetic diversity
- The tip is not accessible \Rightarrow finite reproductive variance



The method of moments Foutel-Rodier, Schertzer '22

The case k = 1: many-to-one $\mathbb{E}_{x} \Big[\sum_{v \in \mathcal{N}_{t}} f(X_{v}) \Big] = \int_{0}^{\infty} f(x, y) p_{t}(x, y)$ with $p_{t}(x, y) \approx e^{\mu x} v_{1}(x) e^{-\mu y} v_{1}(y) =: h(x)\tilde{h}(y)$ with $p_{t}(x, y) \approx e^{\mu x} v_{1}(x) e^{-\mu y} v_{1}(y) =: h(x)\tilde{h}(y)$ reproductive value $\stackrel{\approx}{\longrightarrow} h(x) \int_{0}^{\infty} f(x, y)\tilde{h}(y) dy$ stable configuration

The method of moments Foutel-Rodier, Schertzer '22

The case k = 1: many-to-one \mathcal{N}_t = set of particles alive at t $X_v = \text{position of } v$ $\mathbb{E}_{x}\left[\sum_{v \in \mathcal{N}_{t}} f(X_{v})\right] = \int_{0}^{\infty} f(x, y) p_{t}(x, y)$ with $p_t(x, y) \approx e^{\mu x} v_1(x) e^{-\mu y} v_1(y) =: h(x) \tilde{h}(y)$ The case k = 2 $\mathbb{E}_{x}\left[\sum_{v_{1}\neq v_{2}\in\mathcal{N}_{t}}f(X_{v_{1}})g(X_{v_{2}})\right] = \int_{0}^{\infty}\int_{0}^{t}2r(y)p_{s}(x,y)\left(\int_{0}^{\infty}f(z)p_{t-s}(y,z)dz\right)\left(\int_{0}^{\infty}g(z)p_{t-s}(y,z)dz\right)dsdy$ (0,x) $\approx 2th(x) \Big(\int_0^\infty r(y)\tilde{h}(y) h(y)^2 dy \Big) \Big(\int_0^\infty f(z)\tilde{h}(z) dz \Big) \Big(\int_0^\infty g(z)\tilde{h}(z) dz \Big)$ $-\cdot \Sigma^2$ (s,x)

The method of moments Foutel-Rodier, Schertzer '22





lineages	at the tip	in the bulk	in the bulk	
•	at the tip	at the tip	in the bulk	i

Definition of a **toy model** for **pushed** fronts falling in the universality class of FKPP equation

(1) phase diagram

in accordance with Birzu et al. '20

(2) internal mechanisms

speed of invasion ↔ genetic composition



lineages	at the tip	in the bulk	in the bulk
•	at the tip	at the tip	in the bulk

(2) **analytical criteria** for the emergence of exchangeable genealogies in structured models

spectral theory: fast mixing, reproductive variance

Definition of a **toy model** for **pushed** fronts falling in the universality class of FKPP equation

(1) phase diagram

in accordance with Birzu et al. '20

(2) internal mechanisms

speed of invasion ↔ genetic composition

Tools

(1) extension of the **method of moments** to genealogies of branching diffusions

spinal decomposition

ightarrow location of the ancestral lineages